

Strominger system, Algebroids and Moduli

Carl Tipler

with Roberto Rubio (IMPA) and Mario Garcia-Fernandez (ICMAT)

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Introduction

Candelas-Horowitz-Strominger-Witten (1985) :

Characterization of $N = 1$ supersymmetric Heterotic String compactifications with $M^{10} = M^4 \times X$, with constant dilaton and zero flux $H = 0$, by means of Calabi-Yau manifolds.

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- Thanks to Yau's Theorem, moduli for CY compactifications is identified with complex, Kähler, and bundle moduli;
L. Huang: complex and bundle moduli
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holomorphic Lie algebroid

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- Moduli for non-Kähler heterotic compactifications?

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Recent developments on moduli of non-Kähler heterotic compactifications:

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- De la Ossa-Svanes / Anderson-Gray-Sharpe '14 : inf. moduli related with $H^1(\mathcal{Q})$ for holomorphic double extension

$$\begin{aligned} 0 &\rightarrow T^*X \rightarrow \mathcal{Q} \rightarrow \mathcal{E} \rightarrow 0 \\ 0 &\rightarrow \text{End } V \oplus \text{End } TX \rightarrow \mathcal{E} \rightarrow TX \rightarrow 0 \end{aligned}$$

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Remark: Assume X is $\partial\bar{\partial}$ or $H^{0,2}(X) = 0$

Strominger - Hull geometry

- (X, Ω) complex 3-manifold: X complex with $\Omega \in \Omega_{hol}^{3,0}(X)$
- G : semi-simple compact Lie group
- $P_s \rightarrow X$: principal G -bundle

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- hermitian metric g given by ω (where $\omega = g(J \cdot, \cdot)$),
- A connection (gauge field) on P_s , with curvature F (field strength)
- ∇ unitary connection on (TX, g) , with curvature R

Taking into account the equations of motions:

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Taking into account the equations of motions: **Strominger system**

$$F \wedge \omega^2 = 0, \quad F^{0,2} = 0,$$

$$R \wedge \omega^2 = 0, \quad R^{0,2} = 0,$$

$$d(\|\Omega\|_{\omega} \omega^2) = 0,$$

$$dd^c \omega - \alpha'(\text{tr } R \wedge R - \text{tr } F \wedge F) = 0$$

Remark: as mathematicians, we cut the α' expansion at first order

Parameters and symmetries

M compact, oriented, $6d$ manifold, $P_S \rightarrow M$ a G -bundle. Consider

$$\mathcal{P} = \{(\Omega, \nabla, A, \omega) \in \Omega^3(\mathbb{C}) \times \text{affine connections} \times \text{conn. on } P_S \times \Omega^2 \text{ satisfying (1), (2), (3)}\}$$

- 1 $\Omega \in \Omega^3(\mathbb{C})$ determines an almost complex structure J_Ω
- 2 ω is J_Ω – compatible
- 3 ∇ is a (ω, J_Ω) -unitary connection

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There is a natural **groupoid** of gauge transformations which acts on \mathcal{P} , preserving the solutions

$$\tilde{\mathcal{G}} = \{(g, p) \in \text{Aut}(P_{GI} \times_M P_S) \times \mathcal{P} : g^*(J_\Omega, \omega) = \check{g}^*(J_\Omega, \omega)\}$$

where $p = (\Omega, \nabla, A, \omega) \in \mathcal{P}$, P_{GI} bundle of oriented frames and source/target

$$s(g, p) = p, \quad t(g, p) = g^* p$$

Brute force

Given ρ a solution to ST:

$$\rho = (\Omega, \nabla, A, \omega) \in \mathcal{P}$$

consider the tangent space:

$$T_{\rho}\mathcal{P} \subset \Omega^{3,0+2,1} \oplus \Omega^1(M, \text{End } TM) \oplus \Omega^1(M, \text{ad } P_S) \oplus \Omega^2$$

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Infinitesimal action of the gauge groupoid $\tilde{\mathcal{G}}$:

Brute force

Given p a solution to ST:

$$p = (\Omega, \nabla, A, \omega) \in \mathcal{P}$$

consider the tangent space:

$$T_p \mathcal{P} \subset \Omega^{3,0+2,1} \oplus \Omega^1(M, \text{End } TM) \oplus \Omega^1(M, \text{ad } P_S) \oplus \Omega^2$$

Infinitesimal action of the gauge groupoid Lie $\tilde{\mathcal{G}}$:

$$P: \Omega^0(TM) \oplus \Omega^0(\text{ad } P_S) \oplus \Omega^0(M, \text{End}_{\Omega, \omega} TM) \rightarrow T_p \mathcal{P}$$

$$P(V, \varphi, \psi) = (d\iota_{V^1, 0}\Omega, \iota_V R_\nabla + d^\nabla(\nabla V) + \nabla\psi, \iota_V F_A + d_A\varphi, L_V\omega),$$

Brute force

Linearisation L of ST induces a complex of differential operators

$$0 \rightarrow \text{Lie } \tilde{\mathcal{G}} \xrightarrow{P} T_p \mathcal{P} \xrightarrow{L} \Omega^4(\mathbb{C}) \oplus W \oplus \Omega^5 \oplus \Omega^4$$

with $W = \Omega^{(0,2)+6}(X, \text{End } TX \oplus \text{ad } P_s)$.

Brute force

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(3d arrow $L = \bigoplus_{i=1}^5 L_i$)

$$L_1(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) = d\dot{\Omega}$$

$$L_2(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) = (\bar{\partial}a^{0,1} + \frac{i}{2}F^j, \bar{\partial}\dot{\nabla}^{0,1} + \frac{i}{2}R^j)$$

$$L_3(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) = (d_A \dot{a} \wedge \omega^2 + 2F \wedge \dot{\omega} \wedge \omega, d_{\nabla} \dot{\nabla} \wedge \omega^2 + 2R \wedge \dot{\omega} \wedge \omega)$$

$$L_4(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) = d \left(2\|\Omega_0\|_{\omega_0} \dot{\omega} \wedge \omega_0 + (\|\dot{\Omega}\|_{\omega}) \omega_0^2 \right)$$

$$L_5(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) = \frac{1}{2} d \left(J_0 d\dot{\omega} - J_0(d\omega)^{JJ_0} + 4\alpha' \text{tr}(\dot{\nabla} \wedge R) - 4\alpha' \text{tr}(\dot{a} \wedge F) \right)$$

Space of infinitesimal deformations

Let S^* be the complex

$$0 \longrightarrow S^0 \xrightarrow{P} S^1 \xrightarrow{L} S^2$$

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Proposition (Garcia-Fernandez, Rubio, T.)

S^* is an elliptic complex

Definition:

The finite-dimensional space $H^1(S^*)$ is called space of infinitesimal deformations

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Many questions:

- Link with De la Ossa-Svanes / Anderson-Gray-Sharpe work?
- Decomposition in complex, metric and bundle moduli?
- Integration of infinitesimal deformations, obstructions?

Anomaly VS Bianchi

Strominger system and infinitesimal moduli

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Moduli and AGSOS map

Moduli splitting

- The anomaly equation: local, Green-Schwarz mechanism,

$$H = dB - \alpha'(CS(\nabla) - CS(A))$$

Anomaly VS Bianchi

- The anomaly equation: local, Green-Schwarz mechanism,

$$H = dB - \alpha'(CS(\nabla) - CS(A))$$

- Require *flux quantization*: the local B's glue into a closed 3-form flux, which defines an integral class in $H^3(M, \mathbb{Z})$.

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Recall

$$L_5 = \frac{1}{2} d \left(J_0 d\dot{\omega} - J_0 (d\omega)^{jJ_0} + 4\alpha' \operatorname{tr}(\dot{\nabla} \wedge R) - 4\alpha' \operatorname{tr}(\dot{a} \wedge F) \right).$$

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Define

$$d\mathcal{F} : H^1(S^*) \rightarrow H^3(M, \mathbb{R})$$

$$(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}) \mapsto \frac{1}{2} \left(J_0 d\dot{\omega} - J_0 (d\omega)^{JJ_0} + 4\alpha' \operatorname{tr}(\dot{\nabla} \wedge R) - 4\alpha' \operatorname{tr}(\dot{a} \wedge F) \right).$$

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L_5 is the linearization of Bianchi identity,

$$dd^c\omega - \alpha' (\operatorname{tr} R \wedge R - \operatorname{tr} F \wedge F) = 0,$$

$d\mathcal{F}$ is the linearization of anomaly equation.

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$d\mathcal{F}$ is the linearization of anomaly equation. Set

$$H^1(\tilde{S}^*) := \ker d\mathcal{F},$$

whose elements preserve flux quantization

$\ker d\mathcal{F}$ corresponds to a cohomology group of a complex $\tilde{\mathcal{S}}^*$:

$$\tilde{\mathcal{S}}^0 \subset \mathcal{S}^0 \oplus \Omega^2, \quad \tilde{\mathcal{S}}^1 = \mathcal{S}^1 \oplus \Omega^2.$$

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2-forms play a role of **symmetries**, and play a role as **parameters**.

$\ker d\mathcal{F}$ corresponds to a cohomology group of a complex \tilde{S}^* :

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Substitute L_5 by \tilde{L}_5 (*linearization of anomaly equation*)

$$\tilde{L}_5(\dot{\Omega}, \dot{\nabla}, \dot{a}, \dot{\omega}, b) = db - \frac{1}{2} \left(J_0 d\dot{\omega} - J_0 (d\omega)^{JJ_0} + 4\alpha' \operatorname{tr}(\dot{\nabla} \wedge R) - \dots \right)$$

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The maps

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given by

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define a complex **iff**

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meaning of this equation?

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Equation satisfied by **symmetries of a Courant algebroid** $(E, \langle, \rangle, [,], \pi_T)$
(Baraglia, Rubio, Hitchin) constructed from a solution of the Strominger system

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Infinitesimal symmetries:

$$\text{Lie Aut } E \subset \text{Lie } \tilde{\mathcal{G}} \oplus \Omega^2 \oplus \dots$$

and sub-algebra

$$\text{Lie } \widetilde{\text{Aut } E} \subset \text{Lie Aut } E$$

given by (V, φ, ψ, B) satisfying $(*)$.

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Elements of \tilde{S}^1 : infinitesimal variations of **generalized metrics** $V_+ \subset E$

Exact case $(T + T^*, \langle, \rangle, [,], \pi_T)$

Courant algebroid structure on $T + T^*$:

$$\langle X + \xi, Y + \eta \rangle = \frac{1}{2}(i_X \eta + i_Y \xi), \quad [X + \xi, Y + \eta] = [X, Y] + L_X \eta - i_Y d\xi$$

It has structure group $O(n, n)$, and symmetries include closed 2-forms, B -fields:

$$X + \xi \mapsto X + \xi + i_X B.$$

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Twisted version: an **exact** Courant algebroid $(E, \langle, \rangle, [,], \pi_T)$ (+ axioms)

$$0 \rightarrow T^* \rightarrow E \rightarrow T \rightarrow 0.$$

It is isomorphic, by choosing a (non-canonical) splitting, to

$$(T + T^*, \langle, \rangle, [,]_H := [,] + i_X i_Y H, \pi_T),$$

for some $H \in \Omega_{cl}^3(M)$ (whose class $[H] \in H^3(M)$ parameterizes E).

Transitive case $T^* \rightarrow E \rightarrow T \rightarrow 0$

Given a principal G -bundle P , we obtain by reduction a transitive Courant algebroid E :

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As a vector bundle,

$$E \cong T + \text{ad } P + T^*$$

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Choosing a splitting $T \rightarrow E$, E is isomorphic to

$$(T + \text{ad } P + T^*, \langle, \rangle, [,]_{\theta, H}, \pi_T),$$

where θ is a connection on P (with curvature $F_\theta \in \Omega_{cl}^2(\text{ad } P)$), and $H \in \Omega^3(M)$ such that

$$dH - \langle F_\theta \wedge F_\theta \rangle = 0.$$

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$$(T + \text{ad } P + T^*, \langle \cdot, \cdot \rangle, [\cdot, \cdot]_{\theta, H}, \pi_T),$$

where θ is a connection on P (with curvature $F_\theta \in \Omega_{cl}^2(\text{ad } P)$), and $H \in \Omega^3(M)$ such that

$$dH - \langle F_\theta \wedge F_\theta \rangle = 0.$$

Remark: This last equation is like the Bianchi identity: for any solution to ST, setting

$$P := P_S \times_M P_{GI}$$

build a transitive Courant algebroid using the 3-form H and $F_\theta = F_A + R_\nabla$ the curvature of the product connection $\theta = A \times \nabla$. Choose the pairing so that:

$$\langle F_\theta \wedge F_\theta \rangle = \alpha' (\text{tr } R_\nabla \wedge R_\nabla - \text{tr } F_A \wedge F_A)$$

Strominger system,
Algebroids and
Moduli

Carl Tipler

The bracket

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$$\begin{aligned} [X + r + \xi, Y + t + \eta]_{\theta, H} = & \\ & [X, Y] + L_X \eta - i_Y d\xi + i_Y i_X H \\ & - F_\theta(X, Y) + i_X dt - i_Y dr \\ & + 2c(tdr) + 2c(i_X F_\theta t) - 2c(i_Y F_\theta r). \end{aligned}$$

Generalized metric for E exact:

A metric is a reduction of the frame bundle from $GL(n)$ to $O(n)$.

A generalized metric is a reduction from $O(n, n)$ to $O(n) \times O(n)$.

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$$V_+ \subset E.$$

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A generalized metric



a metric g together with an isotropic splitting $E \cong T + T^*$.

generalized metric for E transitive:

For E transitive, the structure group is $O(t, s)$,

A generalized metric is a reduction to $O(p, q) \times O(t - p, s - q)$.

Admissible metrics (Garcia-Fernandez):

$$V_+ \subset E \text{ and } V_+ \cap T^* = \{0\} \text{ and } \text{rk}(V_+) = \text{rk}(E) - \dim M$$

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Generalized connection

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$$D : \Omega^0(E) \rightarrow \Omega^0(T^* \otimes E),$$

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- Leibniz rule $D_e f e' = \pi(e)(f) e' + f D_e e'$
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Generalized curvature and torsion are defined

The Gualtieri-Bismut connection

Let V_+ be an admissible generalized metric.

We have $V_- := (V_+)^{\perp} \cong T$ and $V_+ \cong E/T^* (\cong T + \text{ad } P)$.

Let $C_+ \cong (\text{ad } P)^{\perp} \subset T + \text{ad } P$.

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Define

$$D_e^B e' := [e_-, e'_+]_+ + [e_+, e'_-]_- + [Ce_-, e'_-]_- + [Ce_+, e'_+]_+,$$

The connection D_B preserves V_{\pm} and has totally skew torsion T_{D_B} .

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But thanks to D^B , we can define a canonical Levi-Civita connection

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Given $\varphi \in C^{\infty}(M)$, D^{LC} modified canonically to D^{φ} , compatible, torsion-free.

Generalized Killing spinor equations

If M is spin, by $V_- \cong T$, introduce the spinor bundle $S_{\pm}(V_-)$.

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If M is spin, by $V_- \cong T$, introduce the spinor bundle $S_{\pm}(V_-)$.

The connection

$$D_{\pm}^{\varphi} : V_- \rightarrow V_- \otimes (V_{\pm})^*,$$

extends to a differential operator on spinors

$$D_{\pm}^{\varphi} : S_{\pm}(V_-) \rightarrow S_{\pm}(V_-) \otimes (V_{\pm})^*,$$

with associated Dirac operator

$$\not{D}_{-}^{\varphi} : S_{+}(V_-) \rightarrow S_{-}(V_-).$$

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Given a generalized metric V_+ , as before, and $\varphi \in C^{\infty}(M)$, the *Killing spinor equations* for a spinor $\eta \in S_{+}(V_-)$ are given by

**Killing spinor equations,
Waldram-Strickland-Constable-Coimbra**

$$D_{+}^{\varphi} \eta = 0,$$

$$\not{D}_{-}^{\varphi} \eta = 0.$$

On a six-dimensional spin-manifold

Theorem (Garcia-Fernandez, Rubio, T.)

Assume that E is exact. Then (V_+, φ, η) is a solution to the Killing spinor equations with $\eta \neq 0$ if and only if $H = 0$, φ is constant and g is a metric with holonomy contained in $SU(3)$.

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Theorem (Garcia-Fernandez, Rubio, T.)

Assume that E is transitive. The Strominger system is equivalent to the Killing spinor equations.

A couple of ideas from the proofs

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$$F_\theta \cdot \eta = 0$$

$$\nabla^- \eta = 0,$$

$$(H - 2d\varphi) \cdot \eta = 0,$$

$$dH - \langle F_\theta \wedge F_\theta \rangle = 0,$$

for $((g, H, \theta), \varphi, \eta)$, where, by $V_- \cong (T, g)$, $\eta \in S_+(T) \cong S_+(V_-)$ (and ∇^- is the Bismut connection with skew-torsion $-H$).

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For the converse in Strominger, given (ω, A, ∇) , one defines $\theta = A \times \nabla$, $H = d^c \omega$ and φ . Note that the Bianchi identity

$$dd^c \omega - (\text{tr } R \wedge R - \text{tr } F_A \wedge F_A) = 0$$

corresponds to

$$dH - \langle F_\theta \wedge F_\theta \rangle = 0.$$

Infinitesimal generalized Killing spinors

Recall

$H^1(S^*) =$ **infinitesimal solutions to ST, modulo symmetries**

and the map (linearisation of anomaly equation):

$$d\mathcal{F} : H^1(S^*) \rightarrow H^3(M, \mathbb{R})$$

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$H^1(\tilde{S}^*) = \text{infinitesimal solutions to ST, compatible with flux quantization, modulo symmetries of } E.$

Infinitesimal generalized Killing spinors

Elements of $\Omega^0(E)$ induce inner symmetries of E . By restriction, we obtain an elliptic complex \widehat{S}^* such that:

$$\begin{array}{ccccccc}
 & & 0 & & & & \\
 & & \downarrow & & & & \\
 & & H^2(M, \mathbb{R}) & & & & \\
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The flux map

Using Kuranishi's technique, can build local moduli spaces \mathcal{M}_{ST} , $\widetilde{\mathcal{M}}_{ST}$ and $\widehat{\mathcal{M}}_{ST}$

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The map

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is the differential of a *Flux map* \mathcal{F} :

$$\mathcal{F} : \mathcal{M}_{ST} \rightarrow H^3(X, \mathbb{R})$$

Flux quantization: restrict to $\mathcal{F}^{-1}(H^3(M, \mathbb{Z}))$

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$d\mathcal{F}$ is a closed $H^3(M, \mathbb{R})$ -valued 1-form

$$d\mathcal{F} \in \Omega^1(\mathcal{M}_{ST}, H^3(M, \mathbb{R})),$$

and provides a foliation (integrating $\text{Ker } d\mathcal{F}$) on the moduli space

The flux map

The leaf of the foliation passing through ρ is the local moduli space

$$\widetilde{\mathcal{M}}_{ST} = \{ \text{solution to gen. Killing spinors eq.} \} / \{ \text{gen. Diffeos} \}$$

$$\widetilde{\mathcal{M}}_{ST} \longrightarrow \mathcal{M}_{ST} \xrightarrow{d\mathcal{F}} H^3(M, \mathbb{R}).$$

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Restrict to *inner symmetries*:

obtain an $H^2(M, \mathbb{R})$ -bundle $\widehat{\mathcal{M}}_{ST}$ over $\widetilde{\mathcal{M}}_{ST}$

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Conjecture

The moduli $\widehat{\mathcal{M}}_{ST}$ carries a natural Kähler structure.

Evidence: there is a natural map $T_\rho \widehat{\mathcal{M}}_{ST} \rightarrow H^1(\mathcal{Q})$ (complex).

AGSOS map

De la Ossa-Svanes/Anderson-Gray-Sharpe interpretation: cocycles in the Dolbeault complex of \mathcal{Q} :

$$\begin{aligned} 0 &\rightarrow T^*X \rightarrow \mathcal{Q} \rightarrow \mathcal{E} \rightarrow 0 \\ 0 &\rightarrow \text{Ad } P_s \oplus \text{End } TX \rightarrow \mathcal{E} \rightarrow TX \rightarrow 0 \end{aligned}$$

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Exact diagram provided X is a $\partial\bar{\partial}$ -manifold or $H^{0,2}(X) = 0$:

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Fu-Yau example (elliptic fibration over $K3$): X is not $\partial\bar{\partial}$ and $h^{0,2}(X) = 1$.

Deformations of pairs

Elliptic complex encoding deformations of pairs of complex structures on M
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$$C^0 \xrightarrow{P_c} C^1 \xrightarrow{L_c} C^2.$$

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Natural questions:

- Image of Ψ ?
- Interpretation of $\text{Ker } \Psi$ as a metric moduli?

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We can define a map:

$$\begin{aligned} \Phi : H^1(S^*) \cap \ker \Psi &\rightarrow H_A^{1,1} \\ [(\dot{\Omega}, \dot{\theta}, \dot{\omega})] &\mapsto [(\dot{\omega} - L_V \omega)^{1,1} - 2c(r^c, F_\theta)] \end{aligned}$$

for some vector field V and $r^c \in \Omega^0(\text{ad } P)$.

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Injectivity of Φ would provide a splitting of infinitesimal moduli as holomorphic components and metric components.